



Lecture 14

Material and Homework



Read:

Chapter 10 Yariv and Yeh

Chapter 11 Yariv and Yeh

Spontaneous Emission Noise

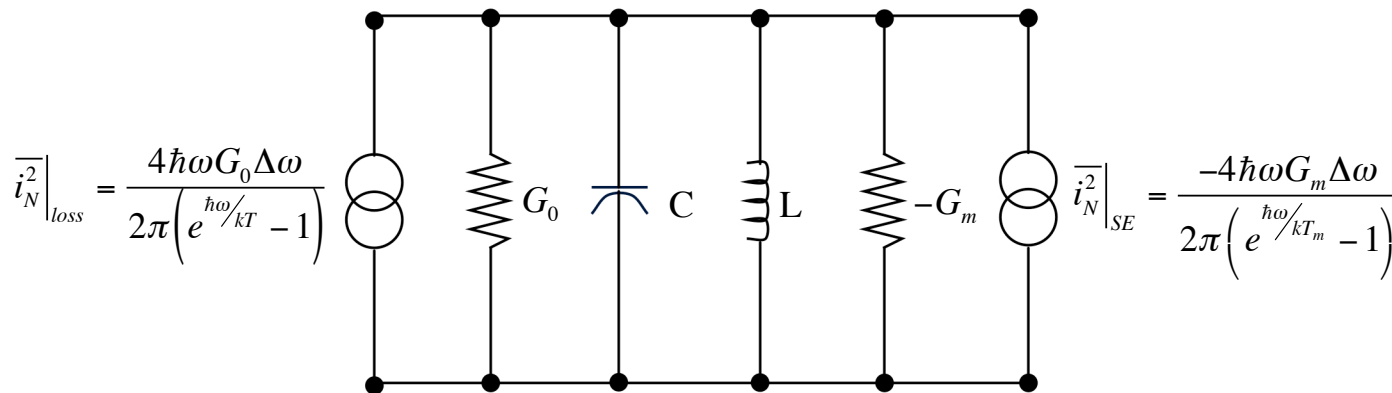


- The spontaneous emission in a laser also contributes to the detected signal to noise ratio (SNR)
 - Photons spontaneously created in the laser cavity will be present at the laser output in addition to stimulated emission.
 - The spontaneous emission and the laser cavity Q will define the laser linewidth

Spontaneous Emission Noise

- We can model the laser cavity as an RLC circuit
 - Conductance $-G_m$ for the gain and G_0 as the loss mechanisms
 - Noise generator $\overline{i_N^2}|_{loss}$ to model Johnson noise due to G_0
 - Noise generator $\overline{i_N^2}|_{SE}$ to model noise due to spontaneous emission for a population inversion represented by temperature T_m is given by $N_2/N_1 = e^{-\hbar\omega/kT_m}$ and $T_m < 0$ for $N_2 > N_1$
 - Using the Boltzman relation

$$\overline{i_N^2}|_{SE} = \frac{-4\hbar\omega G_m \Delta\omega}{2\pi \left(e^{\hbar\omega/kT_m} - 1 \right)} = \frac{4\hbar\omega G_m N_2 \Delta\omega / 2\pi}{N_2 - N_1}$$



Spontaneous Emission Noise

- Defining the quality factor for the equivalent circuit using $\omega_0 = (LC)^{-1/2}$

$$\frac{1}{Q} = \frac{1}{Q_0} - \frac{1}{Q_m} = \frac{G_0}{\sqrt{C/L}} - \frac{G_m}{\sqrt{C/L}} = \frac{G_0}{\omega_0^2 C} - \frac{G_m}{\omega_0^2 C}$$

- The voltage across the equivalent impedance due to the current sources is

$$V(\omega) = \frac{i}{C} \frac{I(\omega)}{\frac{\omega_0^2 - \omega^2}{\omega} + \frac{i\omega_0}{Q}}$$

- The total current can be written assuming the sources are uncorrelated and the total is the sum of the mean squared values

$$\overline{|I(\omega)|^2} = 4\hbar\omega \left(\frac{G_m N_2}{N_2 - N_1} + \frac{G_0}{e^{\hbar\omega/kT} - 1} \right) \frac{d\omega}{2\pi}$$

Spontaneous Emission

- The total power output from the oscillator assuming spontaneous emission dominates Johnson noise (Thermal noise)

$$P = G_0 \int_0^\infty \frac{|V(\omega)|^2}{d\omega} d\omega = \frac{\hbar G_m G_0}{2\pi C^2} \left(\frac{N_2}{N_2 - N_1} \right) \int_0^\infty \frac{\omega d\omega}{(\omega_0 - \omega)^2 + \left(\frac{\omega_0}{2Q} \right)^2} \approx \frac{\hbar G_m G_0 Q}{C^2} \left(\frac{N_2}{N_2 - N_1} \right) \Big|_{\omega \approx \omega_0}$$

- In oscillation, the gain equals losses ($G_0 \sim G_m$) and the power can be written as

$$P = \frac{\hbar G_0^2 Q}{C^2} \left(\frac{N_2}{N_2 - N_1} \right)$$

- Observing that the voltage spectral density near resonance has a linewidth $\Delta\omega = \omega_0/Q$, rewriting using the above equation that relates P to Q, we obtain the Schawlow-Townes linewidth that relates the resonant passive cavity linewidth ($\Delta\nu_{1/2}$) to the laser linewidth, noting that $N_2 = F(P)$

$$\Delta\nu = \frac{\Delta\omega}{2\pi} = \frac{\omega_0}{2\pi P} = \frac{\omega_0}{2\pi Q \frac{\hbar G_0^2}{C^2} \left(\frac{N_2}{N_2 - N_1} \right)} = \frac{2\pi\hbar\nu_0 \left(\Delta\nu_{1/2} \right)^2}{P} \left(\frac{N_2}{N_2 - N_1} \right)$$

Laser Phase Noise



- Laser noise comes from
 - Cavity length fluctuations
 - Vibrational and temperature fluctuations
 - Quantum mechanical fluctuations
- The later, quantum mechanical fluctuations, occurs from the spontaneous emission power being random events that continuously add to the output laser emission field
- Each spontaneous emission event adds a random phase $\Delta\theta$ to the optical (electric) field

$$\Delta\theta = \frac{1}{\sqrt{n}} \cos\phi$$

- After N uncorrelated spontaneous emission events, assuming a uniform distribution of $\Delta\theta$ over $[0, 2\pi]$ the total field phase does an angular random walk, averaging over a large number of spontaneous emission events

$$\langle [\Delta\theta(N)]^2 \rangle = \langle (\Delta\theta_{\text{one emission}})^2 \rangle N = \frac{1}{n} \langle \cos^2 \theta \rangle N$$